Molecular-Scale Force Measurement in a Coiled-Coil Peptide Dimer by Electron Spin Resonance

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Fabrication of multicomponent nanostructures requires the assembly of molecular-scale components into ordered arrays. Biology offers examples of self-assembling structures that form functional entities. Coiled-coil peptides are a particularly interesting class of biological models that naturally form robust multimers, and that can be tuned to yield dimers and trimers as well as large fiber assemblies with predictable morphologies. In order for these structures to find applications as nanodevices new methods are being developed that predict and measure their mechanical properties at the nanoscale level.

In this communication we demonstrate a new experimental method for measuring intercoil forces that is based on electron spin labeling and double electron−electron resonance (DEER) spectroscopy. The model system used for these measurements is derived from the α-helical coiled-coil leucine zipper (LZ) portion (residues 243−281) of the yeast transcriptional activator GCN4 (PDB entry 1YSA), which has recently been characterized by single-site spin labeling. The 4.5 nm × 3 nm leucine zipper subdomain consists of two identical polypeptides.

GCN4-LZ was prepared using solid-phase Fmoc-protection chemistry with a TOAC spin label at residue 248 as shown in Figure 1. Details of the peptide synthesis are given in the Supporting Information. The Multicoil score for the GCN4-LZ sequence (neglecting TOAC) is 0.83, indicating a strong propensity to form a coiled-coil dimer that is reinforced by TOAC’s tendency to adapt a helical backbone conformation.

Distance measurements on nanoscale structures can be obtained by measuring the electron spin−spin dipolar interaction using double electron−electron resonance (DEER) spectroscopy. For the present application, the TOAC spin label was selected because of its rigid fused ring structure (cf. Figure 1), which eliminates motion of the nitroso group relative to the peptide backbone, thus ensuring that spin−spin distance measurements directly reflect the interbackbone distance.

Four-pulse DEER was performed at 65 K as described in the Supporting Information and ref 14, and the results are summarized in Figure 2. The distribution of interspin distances P(r) was obtained from the DEERAnalysis2006 program using model-independent Tikhonov regularization analysis. The nearly ideal shape of the Pake pattern in Figure 2A confirms that orientational selection effects are negligible for the experimental conditions used. The resulting distribution (Figure 2B) showed most of the population at a distance of 2.2 nm, close to the distance of 2.3 nm estimated from a molecular model of the TOAC-labeled dimer based on the GCN4 crystal structure (Figure 1). A small fraction of spins with larger separation is also apparent, which may reflect a minor degree of variation in the backbone structure.

Figure 1. Coiled-coil leucine zipper structure investigated in this work, indicating the structure and position of the TOAC spin label.

Figure 2. (A) Frequency domain DEER signal showing characteristic Pake pattern of an distributed pair of dipoles; (B) solid line shows distribution of distances between spin labels obtained by model-independent Tikhonov analysis of the DEER spectrum. Symbols show distance distribution calculated by the MD-ABF method.
of interaction between coiled-coil units. The distribution of distances in the main fraction (>85%) is very narrow (about 0.14 nm), confirming a compact and well-defined coil structure.

The distance distribution obtained from DEER may be used to calculate the mean force between the halves of the coiled coil. The method is based on the statistical thermodynamic relation between the potential of mean force (PMF) and the probability $P(\xi)$ of finding the system at some particular “reaction coordinate”, $\xi$. The mean force free energy $A(\xi)$ is given by:

$$A(\xi) = -k_B T \ln P(\xi) + A_0$$

where $A_0$ is the standard-state free energy. If one takes $P(\xi)$ to be the spin label distance distribution measured by DEER, one may use this relation to find $A(\xi)$, from which the derivative $-dA(\xi)/d\xi$ may be calculated to give the mean force along $\xi$.

Figure 2B compares the experimental $P(\xi)$ with that calculated from a 4 ns MD trajectory at 298 K using the adaptive biasing force (ABF) method in NAMD$^{20}$ (round symbols). Details of the simulation are given in the Supporting Information. The forces obtained by taking the derivatives of each population distribution curve at 298 K are $110 \pm 10$ pN for the DEER data and $90 \pm 10$ pN for the MD calculation. Given the excellent agreement between the shapes of the calculated and experimental distance distributions, the difference in force may reflect a difference in effective temperatures. That is, if cooling in the DEER sample is not instantaneous, the quenched state of the protein may be closer to the glass transition temperature of the solvent. An effective quenching temperature of 218 K is required to bring the forces into agreement, which is quite close to the glass transition for 30% sucrose. The forces are nevertheless quite comparable to typical protein unfolding forces measured by single-molecule methods.$^{21}$

The method presented here offers a useful complement to existing methods for measuring molecular-scale forces. In contrast to single-molecule methods, it does not require the attachment of the molecule to a macroscopic object such as a bead or atomic force microscope tip. This feature avoids complications from nonspecific binding, permits force measurements under a much wider range of ambient conditions, and also allows one to target significantly smaller structures than are generally accessible to single-molecule measurements. The availability of model-independent algorithms for obtaining distributions$^{16}$ is another advantage of this method, since in principle it can resolve arbitrary distance distributions arising from multiple protein conformations.

On the other hand, the DEER method is inherently an ensemble measurement and therefore cannot directly observe transient events or molecular processivity. It also assumes that the distribution observed under the conditions of DEER (i.e., in frozen solution) reflects the equilibrium distribution of conformations at room temperature. This is valid if the molecular energy surface is sufficiently smooth with a high density of closely spaced states separated by low barriers. In this case, the system is unlikely to be trapped in any particular state in a glassy medium, and the distance distribution obtained from frozen samples should accurately reflect the distribution present at room temperature. The rate of freezing may become important when there is a slow exchange process or global restructuring of the protein.

Accurate force measurements by DEER require a rigidly attached label such as TOAC. The flexible tether in more commonly used spin labels would lead to a broader distance distribution and systematic underestimation of the forces between the labeled protein domains. Although TOAC must generally be incorporated by chemical synthesis, the method described here may be generalized to naturally occurring proteins by taking advantage of other available strategies for immobilizing spin labels relative to the protein.$^{22}$

These results add to our understanding of coiled-coil motifs, which represent an important and common mode of protein–protein interaction. Furthermore, they establish DEER as the only spectroscopic method available for quantitatively measuring the mechanical properties of small peptide-based nanodevices. Such capabilities will be critical for the design of protein-based nanoscale devices with targeted functions.

Acknowledgment. This work was supported by Grants NSF CHE-0443616 and Army W81XWH-06-1-0551 (D.E.B), a NASA NIAC Phase II grant and NSF DMI-0422724 (C.M.), NIH GM60259-01, NSF CHE-0645709, and MRI-0722403 (G.A.L.), and NIH NCRR P41RR16292 (J.H.F.).

Supporting Information Available: Descriptions of the procedures for peptide synthesis, ESR and DEER, and the ABF MD simulation. This material is available free of charge via the Internet at http://pubs.acs.org.

References